ON THE ABSORPTION OF PENETRATING SHOWERS

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SUMMARY

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Several experiments on the absorption in lead of the particles of penetrating showers have been performed under different geometrical conditions. The penetrating showers produced in air are shown to have a high penetrating power whereas those produced invalide locally are strongly absorbed. An interpretation is given by the study of the selection of the penetrating showers by our apparatus. The cross-section for nuclear absorption of the pi-mesons of the locally produced penetrating showers was measured and shown to be smaller than half the geometrical one. Some data on the structure of the PS and on the presence of photons were obtained.

INTRODUCTION

A series of experiments was performed in order to study the absorption of the ionising particles of penetrating showers (PS). We studied the absorption in lead of the paticles of the PS produced in the air as well as those produced in a material located above the arrangement. We call the former atmospheric PS and the latter local PS. The frequency of atmospheric PS is given by the registered frequency without the producing material (background), since accidental coincidences and coincidences due to mesons and their knock-on secondaries are negligible (1) Moreover, the lead shields are thick enough to absorb soft showers (a point that is checked in each experiment) and the position of these shields is such that the registration of a PS produced in them would be very improbable. As we observed that the absorption of the atmospheric PS is very small in general, their absorption in the producing material will be small and the frequency of local PS will be given by the difference between the frequencies with and without the producing material.

Experiment I was performed in Campos do Jordão (844 g/cm²) during the last months of 1949. The other ones were performed in ethers São Paulo (950 g/cm²) in 1950. The G.M. counters used in our experiments, with sensitive area of 1 cm x 15 cm = 130 cm², had glass walls with an external cathode of conducting paper. All statistical errors given below are standard deviations.

EXPERIMENT I

Experiment I is a refinement of an experiment previously described (2). We registered coincidences between four trays of counters (fig.1). Each tray consisted of three counters in parallel.

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Fig.1 shows that at least two particles with sufficient energy to penetrate 19 cm of lead were needed in order to register a coincidence. The PS were absorbed in layers of lead 5 and 10 cm thick placed in position A between the trays of each telescope. We think that the production of PS in this lead is negligible because such a PS would have to cointain at least one energetic particle going backwards. The local PS are produced in 29 g/cm2 of carbon located above the arrangement as shown in fig 1. Our results are collected in table I.We see that the atmospheric PS are not appreciably absorbed by the additional layers of lead, while the polocal production is strongly absorbed in 5 cm of lead. These results are in agreement with those of the preceeding experiment (2). We use in the present experiment carbon as producing material because the atomic weight of carbon is close to the mean atomic weight of the air. The similarity of the results of this experiment and of the preceding one where we used gasoline instead of carbon justifies the indiscriminate use of both materials in our experiments.

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EXPERIMENT II

The purpose of this experiment is to repeat experiment I under different geometrical conditions: we wanted to study the absorption of local PS with an angular aperture sensibly smaller.

The PS were selected by three trays of three counters each (fig.2). The two outer counters in each tray were connected in parallel. We registered sixfold coincidences. For a coincidence to be registered there is to be sixtest at least two particles able to penetrate ten 10 cm of lead. The local PS were produced in carbon.

In the cavities A and B we placed different thicknesses of lead absor-

ber . Our results are collected in table II.

The high background counting rate without absorber is due to the fact that eletrons showers are still registered. From 2,5 cm to 10 cm of lead the background rate shows only a slight variation.

This is due to a small absorption of the atmospheric PS. The locally produced PS are however strongly absorbed, as in experiment I.

EXPERIMENT III

In experiment III the PS are selected by two vertical telescopes, and variable thicknesses of lead are placed above a counter E located between the telescopes (fig 3). We measure in this way the penetrating power of (a) particle of the shower. Each tray of the telescopes consisted of three counters in parallel. Between the trays of each telescope was a permanent layer of 10 cm of lead. We registered both the coincidences among the four trays of the telescopes and the five-fold coincidences among these four trays and counter E.

In a preliminary measurement, without absorber above counter E abouther between the trays of each telescope? we added 5 cm of lead to the screening of the telescopes, and the fourfold coincidence rate of the background decreased sensibly as at shown in table . This suggests that with only the permanent screening recorded of lead some eletrons showers were still registered. When further absorber *** increased the lead screening, the background rate did not changed recorded. appreciably. This shows that no more eletron showers are registered/ with 15 cm of lead and confirms that the absorption of the atmospheric PS is weak. Repeating the same measurements with a producing material (38 g/cm of gasoline) we observed that the locally produced PS/regisrecorded tered by the two telescopes are strongly absorbed (table). This is in agreement with the results of experiments I and II. aborher of 15 con et lead in

Afterwards, with a permanent screening of the telescopes of

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cm of lead and with 67 g/cm² of gasoline of producing material, we placed w various thicknesses of lead above counter E. The fivefold coincidence rates with different thicknesses of lead are shown in table .

The absorber covers only the counter E so that it cannot to either the bank change the fourfold coincidence rate neither of the background nor crowned or to local production of P.S. of the local production. The PS produced in this absorber cannot be registered due to its unfavourable position. These facts are confirmed by our measurements (table 1). This allowed us to collect all of the results of the fourfold coincidences without as well as with the producer. Table xhows does not show any absorption of the fivefold coincidences without producer in agreement with experiments I and II. This allowed us to collect also all of the coincidences due to atmospheric PS. Taking into account that the fivefold coincidences are not statistically independent from the fourfold coincidences, we normalized all our data of fivefold coincidences to the common fourfold coincidence rate and accordingly reduced the statistical errors of the fivefold rates. Applying the method of least squares to the normalized fivefold coincidences rate, under the hypothesis that the absorption follows an expontial law, we found e-1/L =0,982 ± 0,010 where L is the absorption length in cm of lead. The most probable value of L is XXXX 56 cm of lead (640 g/cm2).

We investigated whether or not counter E was struck by photons resulting from the decay of the neutral mesons of the PS which ranks constitute about half of the number of charged mesons (3). As the thickness of the wall of counter E measured in radiation lengths is & 1, this counter has a very low efficiency for the registration of photons. When covered with 1 cm of lead its efficiency

ency will approach unity, in our range of energies. We thus expect the difference between the fivefold coincidence rate with and without 1 cm of lead to yield the frequency of these photons. Table 4 shows that # this frequency is low. We believe that this happens because the angular dispersion of the two photons into which the neutral mesons decay is so large that only a few of them strike counter E. This sets an upper limit of rough #1y 1 Bev to the mean energy of the mesons which go in the direction of counter E, if we assume the number of photons TEXEKXARGHEXAR

Discussion

In our experiments the absorption of atmospheric PS is weak. The atmospheric PS must be produced in a mass of air at least of the order of the mass of producing material used in our experiments because the counting rates are of the same order. By geometrical considerations, we see that the center of this mass of air must be tens of meters above our arrangement. The largest part of the registered atmospheric PS should consist of groups of particles having a very small angular dispersion and in consequence large energy, and should consist mainly of mu-mesons resulting from the decay of pi-mesons. For these reasons these particles should not desapear appreciably by ionization nor suffer nuclear interactions. It is interesting to notice that the predominating factor in the selection of the atmospheric PS is the \$ small angular divergence rather than the thickness of lead used in our recorded experiments so that the registered particles of the atmospheric PS selected by our arrangement have an energy much larger than the minimum energy necessary to cross the lead shielding.

The results of experiments I, II and the absorption of the fourfold coincidences in experiment III, indicate that in the conditions of these experiments the absortion of the locally produced PS is strong, even if we take into account the fact that, to absorb a PS it is sufficient to absorb in only one of the telescopes. The absorption of the particles which strike counter E in the locally produced PS (experiment III) is however weak. Because of this fact and because the cross section for nuclear absorption would have to be larger than the geometrical cross-section we do not believe that the contribution of nuclear interactions to the observed absorption of the fourfold coincidences due to the local PS is important. Thus the principal contribution to the absorption of the PS comes from particles stopped in the lead absorber after having lost their energy by ioni-Eation. Our experimental results are thus that the main part of the particles have an energy just above the amang minimum energy required to cross the lead screening. We believe that the local PS are not selected by the requirement of a particularly small divergenab sor her ce but by the thickness of the lead/screening, which means that the minimum energy of the registered particles of the local PS should be just the energy necessary to penetrate this screening. As the spectral distribution of these particles is believed to fall off rapidly

It is very important to know the criterium by which the PS are selected by a given arrangement. Indeed, the PS are really "penetrating", i.e. their frequency does not depend within large limits on the thickness of the lead screening, only if they are selected by a special requirement such as their small angular divergence. This

with increasing energy, the local PS would be easily absorbed.

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PS may be called "penetrating". The registred local PS, though

of the same nature as the registred atmospheric PS are not really

"penetrating " in the sense given above, for reasons we have already discussed, but they are of course penetrating in the sense that

some of their particles are able to cross large thicknesses of lead.

The small absorption of the local production in the absorber placed about above counter E shows that the largest part of the particles of the PS wich strike this counter have an energy larger than that necessary to penetrate any of the absorbers . This fact indicates that the particles in the center of the PS (which have the direction of the primary particle) of the local PS have in the mean a larger energy than the lateral particled . This is easely understood by the kew laws of conservation of momentum. If ionization does not contribute to the observed absorption, the absorption length observed by us will be due to nuclear absorption and will be 2 400 g/cm2. There is a correction due to the fact that counter E is sometimes struck simultaneously by more than one particle. Taking into account this correction and statistical errors of the mesurements, we are still able to state that the nuclear absorption length cannot be shorter than 350 g/cm2. However because of this xpex the possible contribution of ionization and because of the statistical errors, we cannot give an upper limit to the nuclear absorption length.

Notice that since counter E was not struck by photons, we are measuring the absorption length for charged particles. The measure-

This is not correct. Mesons are freety predomin ant among a particles of un minimum ionizations. This is true not because potous have there are more merous than probans, but because potous have there are appreciably higher than unimimum up to every so much i onization than mentions of the merous where the an initial way to every so much

ments with photographic emulsions on PS of these energies indicate that at least 80% of the charged particles are pi-mesons, while the protons which constitute the remainder could not change our results appreciably, because their absorption length in lead is of the order of 300 g/cm². The distance from the producer to counter E is too small to allow the decay of the pi-mesons initially produced. We conclude that the cross-section for nuclear absorption of charged pi-mesons in lead in our range of energies is less than half the geometrical cross-section. This conclusion is in agreement with the result of other workers (5 - 10).

However our results do not allow us to make a definite statement on the value of the nuclear interaction length of the pi-mesons. Our arrangement is not sensitive to small-angle scattering with a small energy transfer to the lead nucleus. But it is believed that nuclear interactions should produce on the average rather large angle scattering and these were shown by Brown and Mckay (5) to be 4 5 times less frequent that catastrolic interactions. Besides our arrangement has a reasonable efficiency for detecting large angle scattering. Another fact which could mask a possible nuclear interaction is the production by the pi-meson of penetrating secondaries in lead. It is rather difficult to meet thes argument but we may remember that the energy of the particles striking counter E cannot exceed something like 1 Bev. In order to reconcile our results with a geometrical cross-section (this corresponds to a mean range ~ 150 g/cm (for nuclear interaction, as found by the Bristol group (A), more than half of all interactions should result in the production of a PS able to strike counter E

with high efficiency. We do not know if this can happen in our range of rather low energies.

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TABLEI

Experiment I

Hourly rate of fourfold coincidences

Lead thickness in position A	Background (Atmospheric PS)	With carbon	Production in carbon
O cm	1,79±0,16	2,21 = 0,15	+0,42±0,22
5 cm	=1,70=0,13	1,53 ± 0,13	-0,17±0,18
10 cm	1,39 ± 0,10	1,50±0,10	+0,11±0,14

TABLE II

Experiment II

Hourly rate of sixfold coincidences

Thickness of	Background	With	Production	
lead absorbers (Atmospheric		earbon	in carbon (local PS)	
0 cm	1,106±0,069	1,455 = 0,082	~0,7	
2,5 cm	0,656 ±0,056	1,000 ± 0,077	0,344 = 0,095	
5 cm	0,562 ± 0,037	0,728 ± 0,037	0,166 = 0,052	
10 cm	0,443 ± 0,043			

TABLE III

First part of experiment III Hourly rates of coincidences

	thickness		O cm	5 cm	10 cm
Background (atmospheric PS) With gasoline Production in gasoline (local PS)	Background	fourfold rate	1,493 0,083	0,853 0,031	0,978 0,070
	fivefold rate	0,930 0,066	0,518 0,024	0,643 0,057	
	fourfold rate	2,009 0,106	1,175 0,056	1,084 0,090	
	fivefold rate	1,181 0,082	0,672 0,042	0,693 0,071	
	fourfold rate	0,8	0,322 0,064	0,106 0,114	
	gasolina	fivefold rate	0,6	0,154 0,049	0,050 0,091

TABLE IV

Second part of experiment III Hourly rates of coincidences

Lead thickness	s Background		With gasoline	
above counter E	fourfold coincidences	fivefold coincidences	fourfold coincidences	fivefold coincidences
O cm	0,834 ± 0,051	0,518=0,026*	1,301 = 0,074	0,730±0,034++
5 cm	0,805 = 0,058	0,548 = 0,029*	1,195 ± 0,069	0,751 = 0,034 ++
15 em	0,906 = 0,074	0,496 = 0,034*	1,234 ±0,073	0,669 ± 0,036**
20 em	0,911 = 0,080	0,499 = 0,029*	1,186 = 0,064	0,679 + 0,030++
Average rate***	0,853 ± 0,031	0,518 = 0,024	1,226 ± 0,035	
1 cm	0,876±0,081	0,539 = 0,038*	1,198±0,080	0,706±0,040**

^{*} and ** - These statistical errors can be used only when comparing measurements in the same series.

⁻ The measurements with 1 cm were not included when computing the average.